

Actin Depolymerization Drives Actomyosin Ring Contraction during Budding Yeast Cytokinesis

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DOI 10.1016/j.devcel.2012.04.015

SUMMARY

Actin filaments and myosin II are evolutionarily conserved force-generating components of the contractile ring during cytokinesis. Here we show that in budding yeast, actin filament depolymerization plays a major role in actomyosin ring constriction. Cofilin mutation or chemically stabilizing actin filaments attenuate actomyosin ring constriction. Deletion of myosin II motor domain or the myosin regulatory light chain reduced the contraction rate and also the rate of actin depolymerization in the ring. We constructed a quantitative microscopic model of actomyosin ring constriction via filament sliding driven by both actin depolymerization and myosin II motor activity. Model simulations based on experimental measurements support the notion that actin depolymerization is the predominant mechanism for ring constriction. The model predicts invariability of total contraction time regardless of the initial ring size, as originally reported for *C. elegans* embryonic cells. This prediction was validated in yeast cells of different sizes due to different ploidies.

INTRODUCTION

The actomyosin ring is a transient structure built with a network of actin filaments and myosin II motor proteins required for the generation of force for cytokinesis in animal cells and yeast. The actomyosin ring is assembled during specific phases of mitosis and is precisely positioned to bisect the elongating anaphase spindle such that sister chromosomes are segregated to opposing sides of the cleavage plane (Balasubramanian et al., 2004; Barr and Gruneberg, 2007; Field et al., 1999; Pollard, 2010). Cell-cycle kinases activated at mitotic exit signal the initiation of actomyosin ring constriction to drive furrow ingression (Glotzer, 2001; McCollum and Gould, 2001; Pollard, 2010; Wolfe and Gould, 2005; Wolfe et al., 2009; Yoshida et al., 2006). The actomyosin ring disassembles promptly on the completion of

furrow closure. Cytokinesis is a key process in development and growth, and impairment of cytokinesis leads to genome instability (Li, 2007; Normand and King, 2010; Storchova and Pellman, 2004). Considerable progress has been made in recent years to understand the regulatory processes and cytoskeletal components required for actomyosin ring assembly and function (Pollard, 2010). However, the basic mechanism of force generation orchestrated by actin filaments and myosin II has remained elusive.

Inferred from actomyosin force production in striated muscle, it was originally proposed that contractile stress in the actomyosin ring is generated via a “sliding filament” mechanism, in which bipolar myosin motor filaments walk along actin filaments within organized sarcomere-like arrays (Schroeder, 1975). However, both classical and recent studies have noted important distinctions between cytokinetic structures and muscle sarcomeres. Electron microscopy studies in cultured mammalian cells and fission yeast showed that actin filaments are more isotropically oriented with respect to each other than forming strictly antiparallel arrays (Kamasaki et al., 2007; Mabu-chi et al., 1988; Maupin and Pollard, 1986; Sanger and Sanger, 1980; Schroeder, 1973). Furthermore, unlike sarcomeric actomyosin structures, which conserve in mass during cycles of contraction and relaxation, it has been widely observed that actomyosin ring reduces in mass, often manifested as the roughly constant ring width and actin and myosin density as the ring shortens in circumference (Carvalho et al., 2009; Schroeder, 1972; Wu and Pollard, 2005). Consistent with the constant width and density, previous studies reported a constant contraction speed during cleavage furrow ingression (Carvalho et al., 2009; Zumdieck et al., 2007).

Although myosin motor activity is required for cytokinesis in many experimental models, several examples are known where myosin II or myosin II motor activity is not strictly required for cytokinesis (Fang et al., 2010; Gerisch and Weber, 2000; Lord et al., 2005). Particularly relevant to this study, cytokinesis in the budding yeast *Saccharomyces cerevisiae* involves a ring structure composed of actin filaments and myosin II (Myo1), and indeed, in euploid strains of both the S288c and W303a backgrounds (two commonly studied genetic backgrounds), myosin II is essential for cell division (Rancati et al., 2008; Tolliday et al., 2003). Although Myo1 is required for actin assembly in

the ring (Bi et al., 1998), surprisingly, the motor domain of Myo1 is not required for ring constriction or cytokinesis (Fang et al., 2010; Lord et al., 2005). Myosin II-independent cytokinesis has also been well documented in myosin II-deleted *Dictyostelium discoideum* cells (De Lozanne and Spudich, 1987; Knecht and Loomis, 1987; Neujahr et al., 1997; Zang et al., 1997). In this system actin filaments and cortexillin, an actin crosslinker, are thought to be important in addition to contributions from microtubules and relaxation of the polar cortex (Simson et al., 1998; Stock et al., 1999; Weber et al., 1999). Considering the differences between species in the extent by which myosin II is required for cytokinesis, it is unclear whether a common framework can be found to explain the basic force generation mechanism in the actomyosin ring.

A possible conceptual breakthrough is the consideration that, in addition to motor action, filament dynamics in actin bundles may have the ability to generate contractile stress. A recent theoretical paper by Zumdieck et al. (2007) proposed a macroscopic model combining material exchange with force balance. This model was able to recapitulate the observation of roughly constant contraction speed and constant filament density observed in dividing *C. elegans* embryos. The authors also proposed microscopic concepts where parallel or antiparallel actin filament configurations could generate force through either myosin II motor or actin depolymerization in the presence of end-tracking crosslinkers. However, it remained to be tested if actin filament dynamics are indeed crucial for cytokinesis and how such mechanism may be integrated with myosin II motor activity.

In this study, we build upon the work of Zumdieck et al. (2007) to propose a quantitative microscopic model accounting for both motor action and filament dynamics to explain the observed properties of the actomyosin ring during budding yeast cytokinesis. We demonstrate, using a combination of model and experimental analyses, that the primary contractile force during budding yeast cytokinesis results from actin depolymerization mediated by cofilin and myosin II motor activity. Our model also predicts the independence of the contraction time on the initial size of the contractile ring, a phenomenon recently reported in the developing *C. elegans* embryos (Carvalho et al., 2009).

RESULTS

Protein Dynamics and Material Balance during Actomyosin Ring Constriction

We measured the rates describing actin dynamics and ring constriction during cytokinesis first in the wild-type (WT) S288c budding yeast cells. To visualize actin filaments in the contractile ring and due to the inability to directly tag filamentous actin (Doyle and Botstein, 1996), we utilized a probe (iqgCH-GFP), which contains a fragment of Iqg1/Cyk1 (aa 1–446), harboring the actin-binding calponin homology (CH) domain and tagged at the COOH terminus with green fluorescent protein (GFP). The homologous region of Rng2 in fission yeast specifically labels actin filaments in the actomyosin ring, but not other actin structures (Takaine et al., 2009). We confirmed that the iqgCH-GFP also specifically labels contractile ring actin in an F-actin-dependent manner (Figure 1A; see Figure S1C avail-

able online). To observe Myo1 dynamics, a GFP tag was introduced to the 3' end of the *MYO1* ORF at the genomic locus (Fang et al., 2010; Lister et al., 2006) (Figure 1B). Contraction speed observed with these probes was consistent with those previously observed in WT cells (Fang et al., 2010; Lippincott and Li, 1998a; Lister et al., 2006; Lord et al., 2005).

Actin and Myo1 ring constricted at a roughly constant speed (Figure 2A; Movies S1A and S1B). We measured the average fluorescence intensity of iqgCH-GFP or Myo1-GFP in the ring, which reports the relative density of actin filaments or myosin II density, respectively, during ring constriction. The average intensity of the actin probe moderately decreased as the ring constricted, whereas Myo1 density remained roughly constant as previously shown by Lippincott and Li (1998a) (Figures 1 and 2A). The apparent decline in actin density was not due to photobleaching because it occurred only during contraction, and our method of intensity measurement accounts for photobleaching (see Experimental Procedures).

To compute in vivo parameters of actin dynamics from the aforementioned microscopy recording, we started with the material balance equation of Zumdieck et al. (2007):

$$\frac{dc}{dt} = k_p - k_d c - \frac{1}{D} \frac{dD}{dt} c, \quad (1)$$

where k_p and k_d are the rates of actin polymerization and depolymerization, respectively, and D denotes the ring diameter. The last term on the right describes the change in actin density (c) due to contraction in the absence of filament polymerization and depolymerization. Neglecting filament polymerization ($k_p = 0$) during contraction (see below for justification of this assumption), we arrive at the equation

$$\frac{1}{c} \frac{dc}{dt} = -k_d - \frac{1}{D} \frac{dD}{dt}, \quad (2)$$

where all terms in the equation have a dimension of inverse time (s^{-1}). Introducing the normalized rate of contraction $\alpha = -1/D \cdot (dD/dt)$ and the normalized rate of density loss $\beta = -1/c \cdot (dc/dt)$ into Equation 2, we obtained

$$k_d = \alpha + \beta. \quad (3)$$

We note that the relation in Equation 3 is valid only for normalized rates that have the same dimension of s^{-1} . The simple relationship of Equation 3 was used for obtaining actin-depolymerization rate from the experimental estimates of α and β (see Experimental Procedures). Equation 3 also implies that depolymerization contributes to both actin density loss and ring contraction in some proportion, as further explored below.

Effect of Myosin II Motor Inactivation on Ring Constriction and Actin Depolymerization

We next performed measurements to quantify the role of myosin II motor in actomyosin dynamics during cytokinesis. We generated a yeast strain expressing GFP-tagged Myo1 lacking the motor domain (Myo1^{Δm}-GFP) at the native *MYO1* locus. Consistent with the previous finding, the motorless construct rescued the cytokinesis defect of $\Delta myo1$ (Fang et al., 2010), whereas the strain remained euploid (data not shown). However, the contraction rate α was significantly reduced (by 40%)

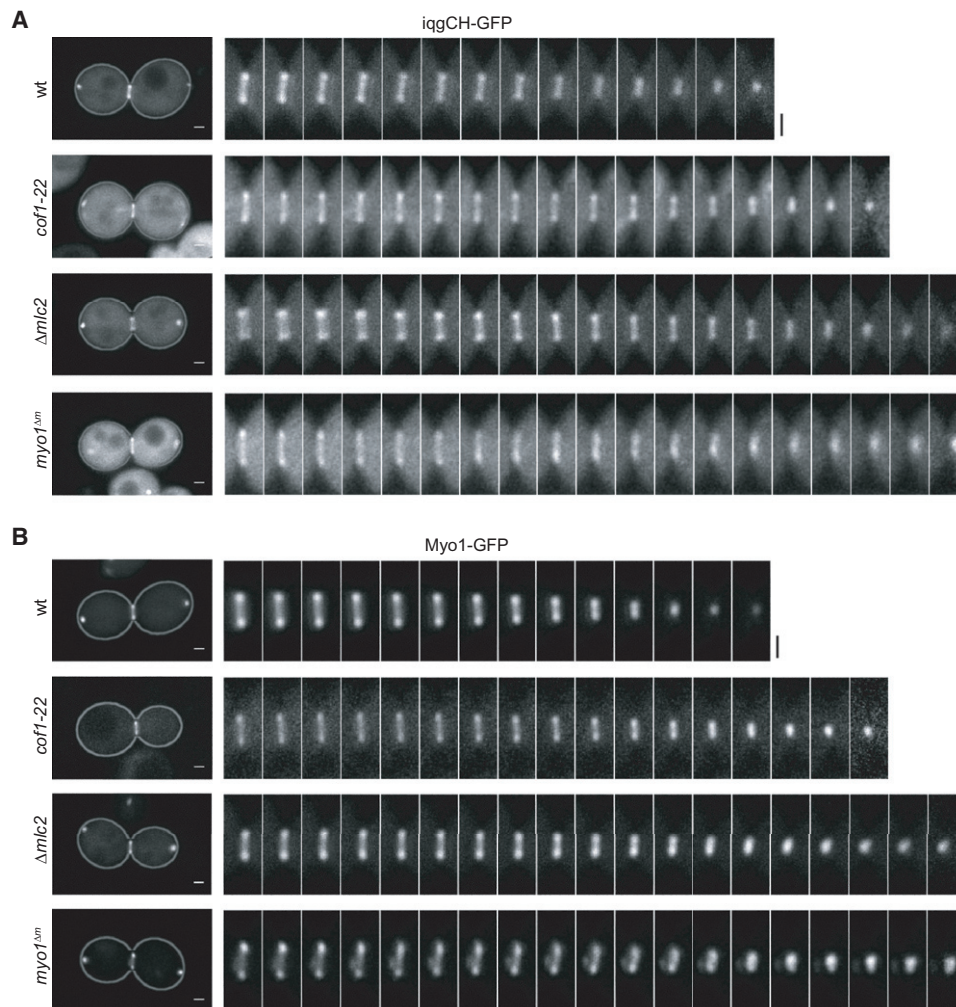


Figure 1. Time-Lapse Imaging of Actin and Myo1 Ring Dynamics in Various Yeast Strains

(A) Wild-type haploid (WT), *cof1-22*, $\Delta mlc2$, and *myo1 Δm* cells expressing the actin probe iqqCH-GFP and Spc42-mCherry-spindle pole body probe (seen as the two dots near the cell poles).

(B) WT, *cof1-22*, $\Delta mlc2$, and *myo1 Δm* cells expressing Myo1-GFP and Spc42-mCherry. Shown are representative time-lapse images of constricting rings for the actin probe and Myo1-GFP. The images of each montage are 23.6 s apart, and each image represents the temporal average of six consecutive frames (3.93 s apart) in the corresponding movie. The first panel in each row shows the whole-cell image corresponding to the first image of the montage to the right that corresponds to a cropped area around the bud neck. Scale bars are 1 μm .

See also [Movies S1](#) and [S2](#).

compared to the WT rate, and the actin density loss rate was also slightly reduced (Figure 2B; [Movies S2E](#) and [S2F](#); [Table 1](#)). k_d estimate based on [Equation 3](#) found *myo1 Δm* to reduce the actin-depolymerization rate by 42%. Mlc2 is the regulatory light chain for the yeast myosin II expected to be important for Myo1 motor activity (Luo et al., 2004). We tested if $\Delta mlc2$ exhibited a similar effect. Indeed, $\Delta mlc2$ also significantly reduced the contraction rate and the rate of actin depolymerization (Figures 1 and 2B; [Movies S2C](#) and [S2D](#); [Table 1](#)). However, unlike the case in fission yeast (see [YFP-MYO2](#) experiments in [Stark et al., 2010](#)), an extra copy of *MYO1-GFP* in budding yeast merely broadened the ring along the bud-mother axis without any increase in Myo1 concentration or alteration of ring contraction dynamics (data not shown). The aforementioned results nonetheless suggest that yeast myosin II motor activity is

important for actomyosin ring constriction and actin depolymerization during the constriction process.

Evidence that Actin Depolymerization Is Important for Actomyosin Ring Constriction

To evaluate the role for actin depolymerization in cytokinesis, we tested the effect of blocking actin depolymerization by the actin-stabilizing drug, jasplakinolide ([Ayscough, 2000](#); [Ayscough et al., 1997](#); [Bubb et al., 1994, 2000](#); [Cramer, 1999](#); [Lee et al., 1998](#)), on actomyosin ring constriction. First, a titration experiment was performed to determine a concentration of jasplakinolide sufficient to cause a considerable reduction in the rate of actin disassembly in yeast cells without inducing gross disorganization of actin structures (Figure S1D; [Movie S3B](#)). We then imaged actomyosin ring constriction in the presence of this

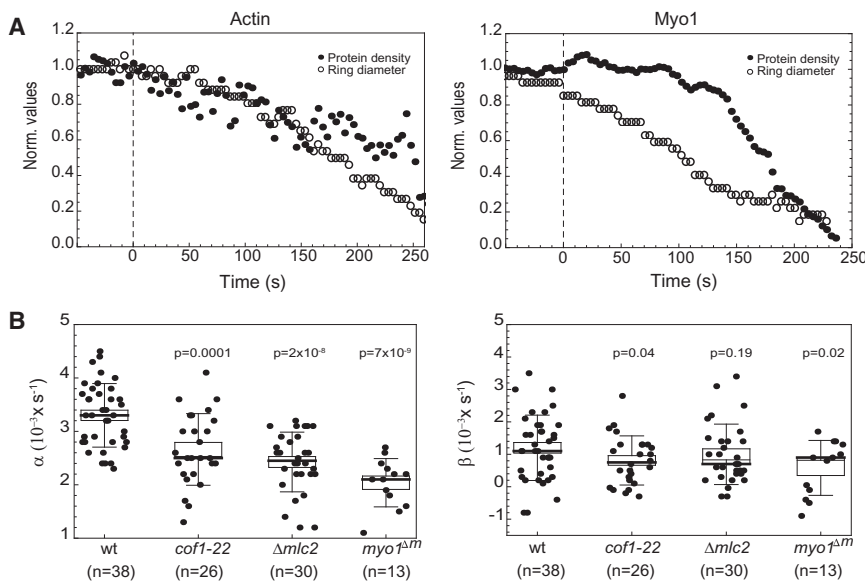


Figure 2. Quantification of the Parameters of Actin Dynamics in Various Yeast Strains

(A) Kinetic profile ring diameter (open circles) and protein density (closed circles) change over time during contraction observed with the actin probe iqqCH-GFP and Myo1-GFP.

(B) Actin ring contraction rate (α) and density loss rates (β) for WT, *cof1-22*, $\Delta mlc2$, and *myo1 Δm* . Box range represents the SEM, whiskers are the SD, the small square is the mean, and the line inside the box is the median. The p values above each plot are for comparison with WT.

See also Figures S1A and S1B.

The Synergy between Motor Action and Actin Depolymerization in Ring Constriction

The results above revealed two molecular functions contributing to actomyosin ring contraction in budding yeast: one relying on actin depolymerization, the other requiring the myosin II motor. These func-

tion concentration (100 μ M) of jasplakinolide. The jasplakinolide treatment led to a significant decrease (26%) in the contraction rate compared to DMSO-treated cells (Figures 3A, 4A, and 4B; Movie S3; Table 1). Interestingly, the drug had a minimal and insignificant effect on actin density change in the contractile ring compared to the DMSO-treated control (Figure 4C), suggesting that the actin depolymerization blocked by jasplakinolide (by 21% compared to DMSO control) was quantitatively correlated with the slowdown of ring constriction.

To further investigate the role of actin depolymerization in ring constriction, we tested the effects of a mutation in *COF1*, encoding cofilin, the main actin-depolymerizing factor in yeast (Lappalainen and Drubin, 1997; Moon et al., 1993). *cof1-22* was previously shown to impair the actin-depolymerization activity of Cof1 in vivo and in vitro (Lappalainen and Drubin, 1997). *cof1-22* cells are able to assemble a morphologically normal actomyosin ring. Consistent with the effect of jasplakinolide, the rate of ring constriction was significantly reduced, whereas the rate of actin density change in the ring was not significantly altered (Figures 1 and 2B; Movie S1; Table 1). Based on Equation 3, *cof1-22* reduced actin-depolymerization rate by 25% compared to the WT rate, consistent with the moderate effect of the mutant protein on actin depolymerization observed in vitro (Lappalainen and Drubin, 1997). However, this reduction of actin-depolymerization rate was accompanied with a similar level of reduction (by 22%) in ring constriction rate (Figure 2B; Table 1), suggesting that the cofilin-mediated actin depolymerization largely contributed to actomyosin ring constriction. We then treated *cof1-22* mutant cells with 100 μ M jasplakinolide (to combine the effects of both methods for blocking actin depolymerization). Under this condition, in 6 of the 35 cells imaged, actomyosin ring contraction was completely blocked, whereas in the remaining 29 the ring contracted with a drastically reduced rate (by 65% compared to WT DMSO control) (Figures 3C and 4; Table 1; Movie S5). These results strongly suggest that actin depolymerization plays a critical role in actomyosin ring contraction.

tions are not independent because the latter also affects actin depolymerization during constriction. To examine the synergy between depolymerization and motor action, we tested the effect of 100 μ M jasplakinolide on the constriction of Myo1 Δm ring. In 16 out of 32 cells examined, jasplakinolide completely blocked Myo1 Δm ring contraction (Figure 3B; Movie S4C), and in the remaining 16 the contraction rate as well as the actin-depolymerization rate were severely reduced (Figures 3B and 4; Movies S4A and S4B; Table 1). Note that in all the cells imaged, in the absence or presence of the very slow contraction, no buildup of actin density was observed. Treatment of these cells with LatA did not lead to reduced actin density in the ring (Figure S1E), suggesting that the stable actin density was not due to balanced polymerization and depolymerization. This observation supports the assumption that actin polymerization was negligible during the time of ring constriction. Consistent with this assumption for actin, FRAP experiment suggests that Myo1 does not undergo turnover during the constriction phase (Figure S1F).

A Quantitative Model of Actomyosin Ring Contraction

To build a mechanistic framework explaining the experimental observations described above and to allow quantitative predictions of actomyosin ring behavior in the cell, we constructed a bottom-up model for ring contraction starting from a set of microscopic elements and their interactions and computationally predicted macroscopic variables such as the rates of contraction and actin density loss in the actomyosin ring. The microscopic elements and our basic assumptions of their interactions are as follows:

- (1) The contractile ring is made up of actin filaments, bipolar myosin II motors, actin-depolymerizing protein cofilin, and an actin crosslinker with the possibility to maintain association with shortening filament ends (Figures 5A and S2A).
- (2) Based on the classical assumption of myosin II motor function, Myo1 can slide an actin filament relative to

Table 1. Experimental and Model Predicted Parameters of Actin Ring Dynamics

	No.	α		β		k_d
		Experimental	Predicted	Experimental	Predicted	
WT	38	3.3 ± 0.1	3.3 ± 0.04	1.2 ± 0.2	1.5 ± 0.06	4.5
WT + DMSO	14	3.5 ± 0.2	3.5 ± 0.05	1.3 ± 0.4	1.6 ± 0.09	4.8
WT + Jaspl 100 μ M	17	2.6 ± 0.1	2.8 ± 0.05	1.2 ± 0.2	1.3 ± 0.08	3.8
<i>cof1-22</i>	26	2.6 ± 0.1	2.6 ± 0.04	0.8 ± 0.2	0.9 ± 0.07	3.4
<i>cof1-22</i> + DMSO	45	3.1 ± 0.2	3.4 ± 0.04	1.3 ± 0.2	0.9 ± 0.01	4.4
<i>cof1-22</i> + Jaspl 100 μ M	29	1.1 ± 0.1	1.2 ± 0.04	0.8 ± 0.1	0.8 ± 0.01	1.9
<i>cof1-22</i> + Jaspl 100 μ M	6	0	0	0	0	0
Δ <i>mlc2</i>	30	2.4 ± 0.1	2.6 ± 0.04	1.1 ± 0.2	1.2 ± 0.06	3.5
<i>myo1Δ^m</i>	13	2.0 ± 0.1	1.9 ± 0.03	0.6 ± 0.2	0.8 ± 0.05	2.6
<i>myo1Δ^m</i> + DMSO	9	2.0 ± 0.1	1.9 ± 0.03	0.5 ± 0.1	1.0 ± 0.04	2.6
<i>myo1Δ^m</i> + Jaspl 100 μ M	16	1.6 ± 0.1	1.4 ± 0.02	0.3 ± 0.2	0.6 ± 0.03	1.9
<i>myo1Δ^m</i> + Jaspl 100 μ M	16	0	0	0	0	0

The model-predicted values were averages from 100 simulations. K_d values are experimental and used in model simulations. The error estimates were SEM. No., number of cells imaged for the experimental values.

the other in a given filament pair at a single motor rate of v (in unit of s^{-1}).

- (3) A second type of actin sliding can be induced by actin depolymerization, preferentially at or near the pointed end, combined with an actin crosslinker (Zumdieck et al., 2007) (Figure 5A). We introduce the rate k (in unit of s^{-1}) for filament sliding caused by actin depolymerization, and parameter γ , representing the average fraction of time by which a crosslinker is in position to reattach to the new filament end after a depolymerization event, such as a cut by cofilin (Figure 5A, left column). Thus, γ is essentially the frequency of productive crosslinking to allow depolymerization to cause filament sliding (Figure 5A).
- (4) The rate of the motor v is expected to be lower than the rate of depolymerization k because otherwise there would be a general buildup of density in the ring during contraction, which was not normally observed (Figure 1B). We further estimate v to be roughly $k/4$ (Supplemental Experimental Procedures, where we also discuss the estimates for v and γ).
- (5) For pairs of filaments we consider all possible configurations in regard to filament orientations and positions of the motor and crosslinker. The formula for the combined rate of sliding of one filament (the lower one in each entry of Figure S2A) with respect to the other (upper) is provided for each case (Figure S2A). Examples are also shown in Figure 5C.

Having specified the structural components, we consider a relatively simple contractile ring represented by a one-dimensional chain of the number (n) of actin filaments arranged in a ring (Figure 5B; Movie S6). The initial mean size of each filament is L_0 with the SD δL ; the mean size of the overlapping region of the adjacent filaments is S_0 with the SD δS . Addition of each filament in the computation construction of the ring in general leads to lengthening of the structure. The initial orientations of the filaments are generated by a random

process controlled by a clustering coefficient ζ describing filament clustering (grouping of parallel filaments, Supplemental Experimental Procedures). The extreme values of ζ are +1 and -1 and correspond to ordered ring structures, namely, all parallel filaments and filaments with alternating orientations, respectively. The values of ζ far from these extremes correspond to random filament orientations. A key simplification making the computation feasible is that the i^{th} filament interacts only with two adjacent $(i-1)^{\text{th}}$ and $(i+1)^{\text{th}}$ filaments. After a given time step, we compute filament positions as a result of sliding based on the formulas described in Figure S2A for any given filament configuration. Thus, Figure S2A was essentially used as a lookup table for each step of the computation (see Supplemental Experimental Procedures for additional details).

Consider several examples of the displacement rate computation using Figure S2A. The simplest structure is shown as the first example in Figure 5C (top pair). The relative filament positions do not allow any activity of the crosslinkers at the pointed end (tail of the arrows in the illustration), so that only the myosin motor activity can cause the lower filament to move to the right (positive direction) with the speed equal to $2v$. A more complex configuration is shown as the structure corresponding to the middle example in Figure 5C, where a myosin motor and a single crosslinker slide the lower filament in same direction relative to the top filament. During the time step Δt , the crosslinker is active for $\gamma \Delta t$ with the rate k , whereas the myosin motor contributes to the displacement for $(1-\gamma)\Delta t$ with the speed v that leads to a shift of the lower filament to the left with the rate $\gamma k + (1-\gamma)v$. An even more complex case is the bottom example in Figure 5C showing two crosslinkers competing with a myosin motor. The probability of one crosslinker to be inactive is $(1-\gamma)$; assuming that they are working independently, we find total probability of both crosslinkers to be inactive as $(1-\gamma)^2$ and thus the motor-driven displacement rate equal to $(1-\gamma)^2 v$. The contribution of the crosslinkers can be found given that at least one of them is active $1 - (1-\gamma)^2 = (2\gamma - \gamma^2)$ fraction of time, leading to a rate of $(2\gamma - \gamma^2)k$.

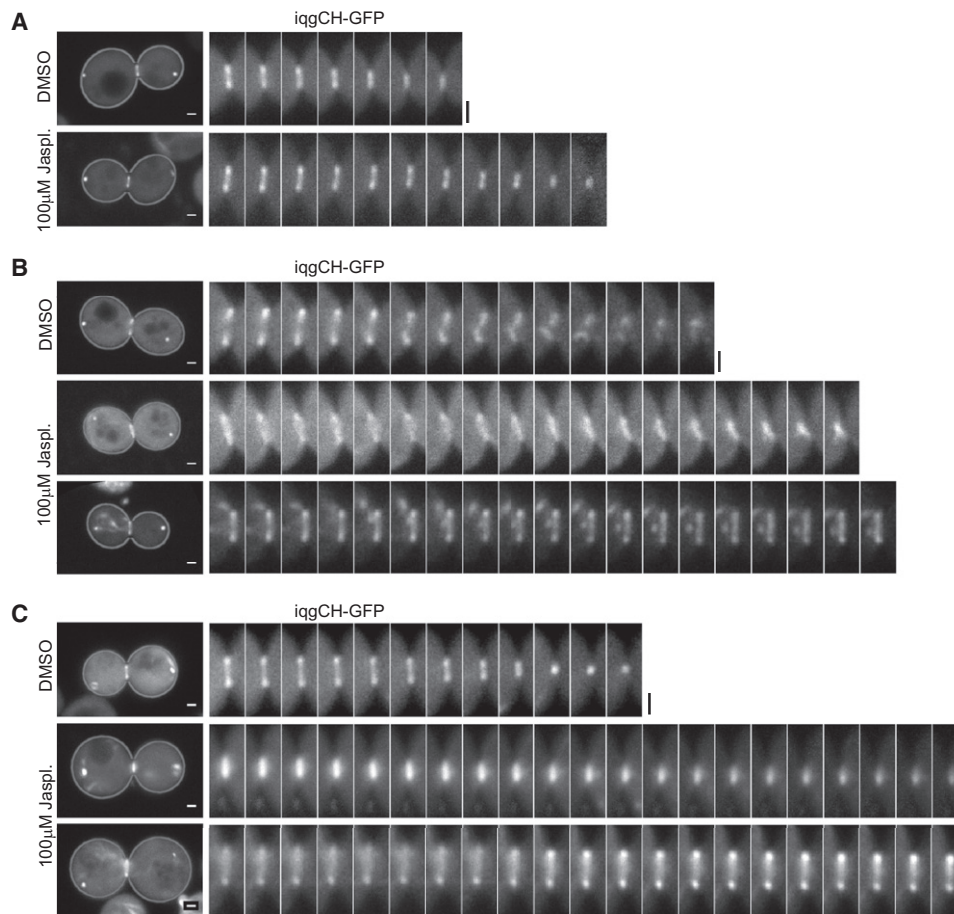


Figure 3. Inhibition of Actin Depolymerization with Jasplakinolide Slows Down Contraction

(A) Representative time-lapse sequences of constricting rings in WT cells expressing the actin probe iqqCH-GFP and Spc42-mCherry and treated with DMSO or 100 μ M of jasplakinolide (Jaspl.).

(B) Representative time-lapse sequences of constricting contractile rings in *myo1 Δ m* cell expressing the actin probe iqqCH-GFP and Spc42-mCherry and treated with DMSO or 100 μ M of jasplakinolide (bottom two rows). The top montage for the jasplakinolide-treated cells shows a slow contracting ring, representing 16 of 32 cells observed in this experiment; the bottom montage shows a noncontracting ring, representing also 16 of 32 cells observed (see [Movie S4](#) to observe the movement of the SPBs indicating mitotic exit).

(C) Representative time-lapse sequences of constricting contractile rings in *cof1-22* cell expressing the actin probe iqqCH-GFP and Spc42-mCherry and treated with DMSO or 100 μ M of jasplakinolide (bottom two rows). The top montage for the jasplakinolide-treated cells shows a slow contracting ring, representing 29 of 35 cells observed in this experiment; the bottom montage shows a noncontracting ring, representing also 6 of 35 cells observed (see [Movie S5](#) to observe the movement of the SPBs indicating mitotic exit). The images of each montage are 39.3 s apart, and each image represents the temporal average of ten consecutive frames (3.93 s apart) in the corresponding movie. The first panel in each row shows the whole-cell image corresponding to the first image of the montage to the right that corresponds to a cropped area around the bud neck. Scale bars are 1 μ m. See also [Movies S3](#), [S4](#), and [S5](#).

Thus, the total rate of displacement of the lower filament to the left is $(1 - \gamma)^2 v + (2\gamma - \gamma^2)k$.

At the end of each time step of a simulation, the length of each filament, the size of the entire filament structure, and the current value of the actin density in the contracting ring were computed. This procedure was repeated until a certain minimal size of the ring was reached. The resulting data allowed us to compute the normalized rate of contraction α by fitting the data to the linear function $L(1 - \alpha t)$, where L denotes the initial size of the filament structure. Similarly, the normalized actin density loss rate β was found by fitting of the computed actin density changes to the function $C(1 - \beta t)$, where C is the initial actin density in the ring, which can be estimated as $C \sim L_0 / (L_0 - S_0)$. A real

contractile ring can be considered as being composed of many simple contractile rings as modeled above, linked laterally with one another. Thus, the real contraction and density loss rates would be equal to the average rate of contraction for a number of such structures, built and contracting with the same parameters.

To search for reasonable parameter values of γ and ζ , we used the experimental data ([Table 1](#)) for two characteristic cases, WT and Myo1 motor deletion, to compute all possible values of γ and ζ that produce a α/β ratio in the range of 2.0–4.0 as observed. The overlap of the parameter spaces from these two cases gave possible values of γ and ζ ([Figure 5D](#)). γ is in the range of 0.25–0.35 that is similar to the estimate given in [Supplemental](#)

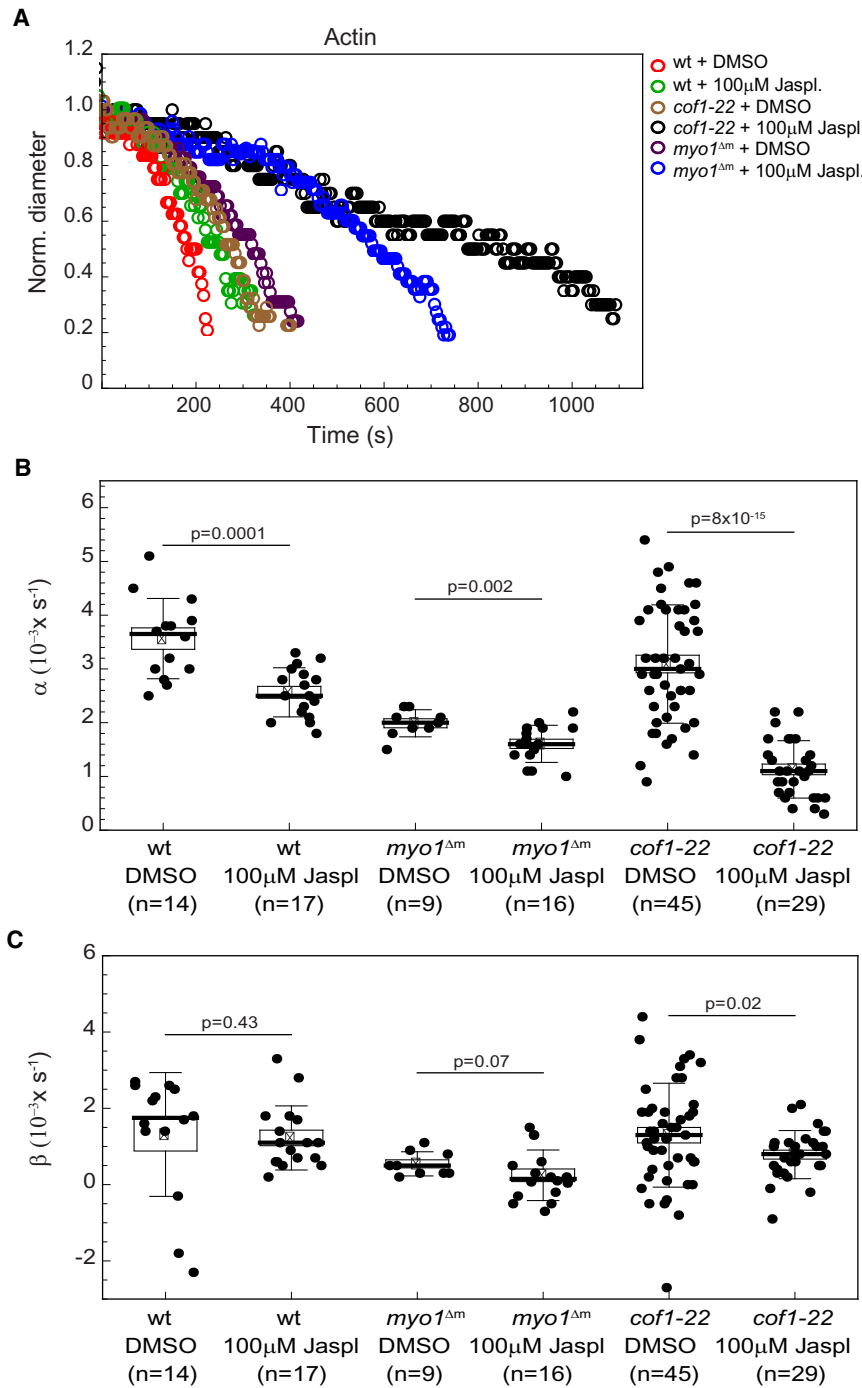


Figure 4. Quantification of Actin Dynamics in the Contractile Ring in the Presence of Jasplakinolide

(A) Representative plots of normalized (Norm.) actin ring diameter versus time for WT, *myo1 Δm* , and *cof1-22* cells treated with DMSO or 100 μ M of jasplakinolide (Jaspl.). (B and C) Contraction rate α (B) and protein density loss rate β (C) for WT, *myo1 Δm* , and *cof1-22* cells treated with DMSO or 100 μ M of jasplakinolide. Box plots are as described in Figure 2B.

constant contraction speed, consistent with the experimental observation (compare Figure 5E). The predicted values of α and β , which were nondimensional, when scaled with the contraction rate α observed for WT cells, were remarkably consistent with the experimentally measured ones under all conditions (Figures 5F and 5G; Table 1). The consistency between measured versus predicted α and β values was also observed when γ and ζ pairs near the edges of the allowable parameter space were used (Figure S2B).

It is important to note that our model rules allow both local contraction and expansion on the level of individual filament pairs (Figure 6). The overall contraction ($\alpha > 0$) or expansion ($\alpha < 0$) rate for an arbitrary ring structure can be estimated from averaging all the local contraction/expansion events for individual filament pairs. To examine the effects of various parameters on overall ring contraction or expansion, we estimated α with a large number of random ring structures. The result of this analysis shows positive contraction rate α for a wide range of the filament orientation parameter ζ in the presence of actin depolymerization ($k > 0$) (Figures 6A and 6C). When the rate of actin depolymerization is zero, the myosin motor activity alone can lead to either contraction or expansion depending on filament orientation and the frequency

of crosslinking for the filament configurations shown in Figure S2 (Figure 6B).

A Theoretical Prediction of the Mechanism by which Myosin II Drives Ring Constriction

The observation that deletion of the Myo1 motor domain or *MLC2* reduced the rate of actin depolymerization suggested that myosin II could contribute to ring constriction through two possible mechanisms: (1) the classical motor-driven filament sliding, and (2) actin depolymerization. We used our model

Experimental Procedures, meaning that less than 40% of the depolymerization events need to produce filament sliding. ζ is in the range of 0.2–0.4, representing filament orientations where there are slightly more parallel filaments than random. A direct validation of the model is to use parameter values from different regions of the parameter space to predict the contraction rate α and actin density loss rate β for all cases where these parameters were measured experimentally, given the estimated k_d values. Using γ and ζ values of 0.27 and 0.37, respectively, the computed dynamics of ring contraction in WT showed roughly

of crosslinking for the filament configurations shown in Figure S2 (Figure 6B).

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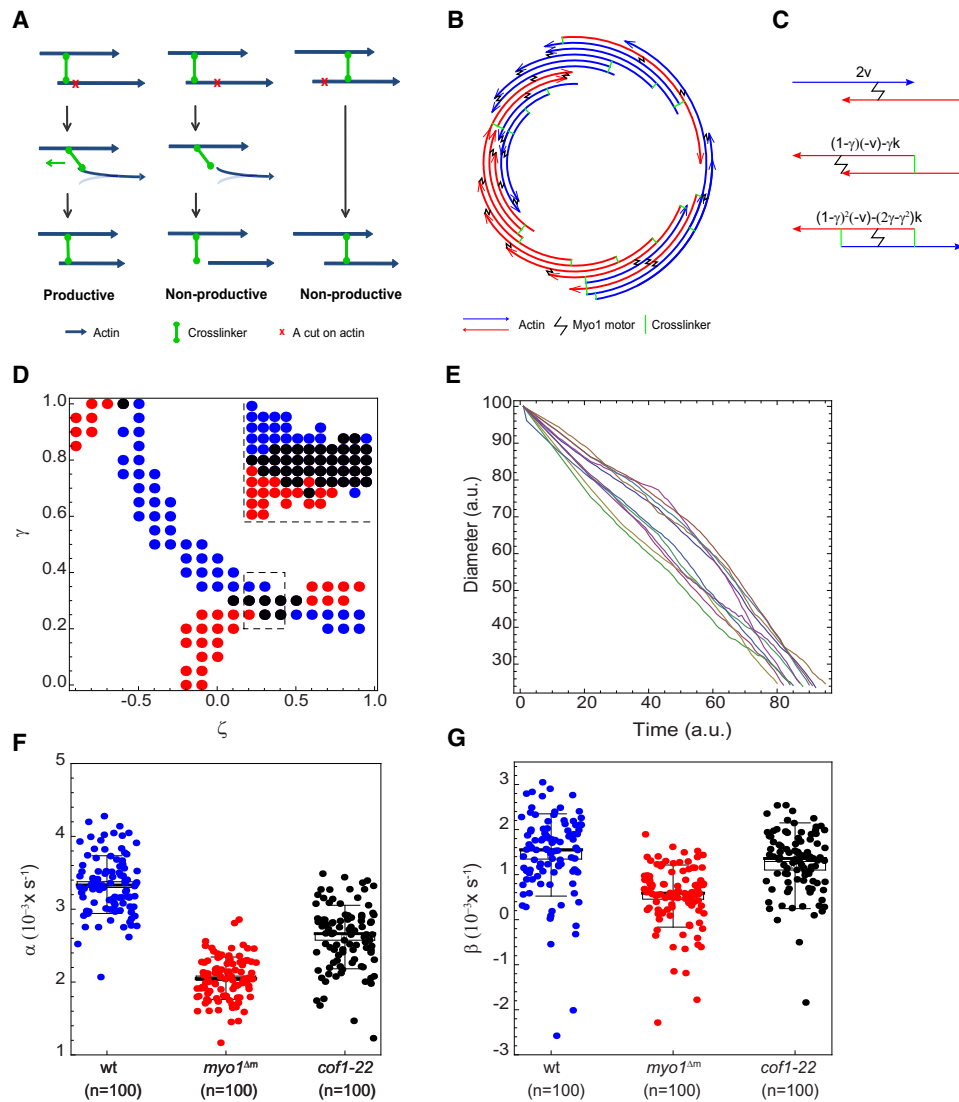


Figure 5. A Quantitative Microscopic Model for the Contraction of the Yeast Actomyosin Ring

(A) Schematic diagram illustrating an example of productive or nonproductive actin-depolymerization event (e.g., a cut produced by cofilin) for generating sliding between two actin filaments (long blue arrows). In the former case (left column), thermal fluctuation of the crosslinker and/or actin filament allows the crosslinker to reattach to the newly produced actin end, and the resulting elastic force could lead to sliding of the lower filament (in the direction marked by the small green arrow) relative to the upper one. In the latter case, actin depolymerization is nonproductive in generating filament sliding possibly due to the cut being too far away from the crosslinker (center column) or being between the filament end and the crosslinker (right column). A larger value for the model parameter γ implies that the former scenario occurs at a higher frequency relative to the latter.

(B) An example of a contractile ring constructed based on the assumptions of the model showing myosin motors (black) and crosslinkers (green) with respect to Figure S2A. Colored arrows show the clockwise (red) and anticlockwise (blue) filament orientations. See also Movie S6.

(C) Examples of basic filament structures showing myosin motors and crosslinkers in the productive positions and giving formula for displacement rate of the lower filament relative to the upper one. See also Figure S2A.

(D) Parameter space for the clustering coefficient ζ and crosslinkers activity γ values for which the ratio of the contraction rate α to the actin loss rate β is in the experimentally observed interval $2 < \alpha/\beta < 4$ for the haploid WT (blue dots) and *myo1^{Δm}* (red dots) cells. The intersection of the parameter space corresponding to these two genotypes is shown in black. Each dot represents 200 simulations for randomly generated initial filament structures. The inset zooms into the rectangular area bounded by the broken lines and shows similar results in higher resolution (with the step size of 0.02 in both parameters); each dot corresponds to 300 simulations.

(E) Representative simulated contraction curves normalized to the initial ring diameter show roughly linear diameter decrease over time, similar to experimentally observed kinetic profile of contractions (compare with Figure 2A).

(F and G) Simulation of ring contraction (F) and actin density loss (G) rates for $\zeta = 0.37$ and $\gamma = 0.27$ for the WT, *myo1^{Δm}*, and *cof1-22* mutant cells successfully reproduces the contraction rates as experimentally observed (compare to Figure 2B). All computed nondimensional rates were scaled to the corresponding mean experimental contraction rate α for WT. Box plots are as described in Figure 2B.

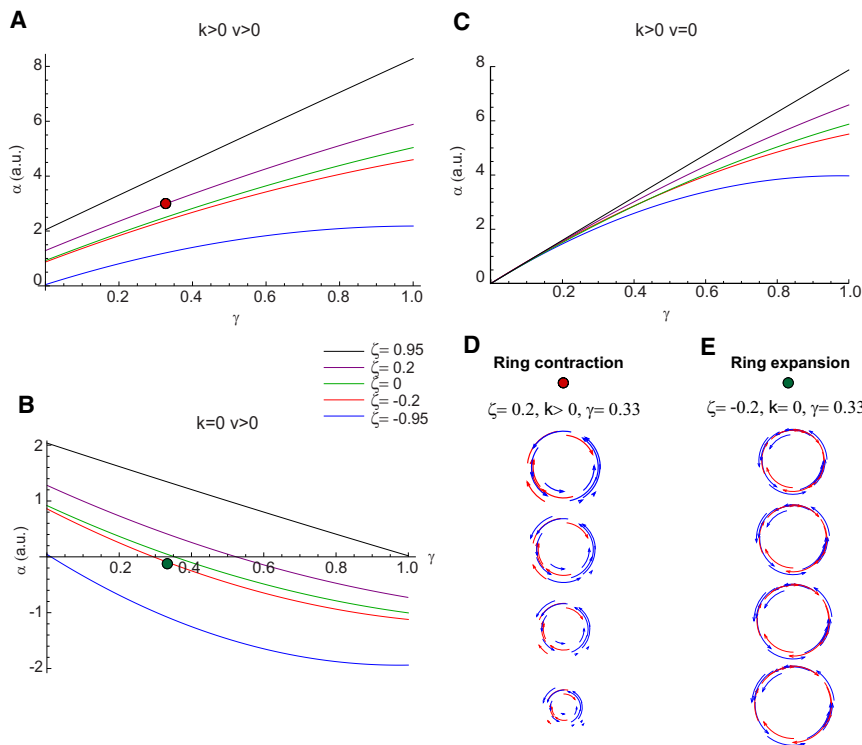


Figure 6. Parameter Analysis of Contraction versus Expansion

Comparison of the contraction rate estimates for different values of the filament orientation factor ζ and crosslinking frequency γ in three cases: both contraction mechanisms are active (A), only myosin motor-based contraction mechanism is active (B), and only depolymerization-driven mechanism is active (C).

(A) The ring contracts (positive α) with a wide range of filament orientation when both myosin motor and depolymerization-based mechanisms are active ($k > 0$, $v > 0$). An example of the simulated ring dynamics is shown in (D). The red dot in (A) corresponds to the parameter values used in model ring simulation in (D).

(B) When filament sliding is driven solely by myosin motor, the ring can either contract (positive α) or expand (negative α) depending on relative filament orientation and crosslinking efficiency. The green dot corresponds to the parameter values used in model ring simulation in (E).

(C) When filament sliding is driven solely with depolymerization-based mechanism, the ring always contracts when γ is nonzero.

(D) An example of a simulated contracting ring corresponding to the red dot in (A).

(E) An example of a simulated expanding ring corresponding to the green dot in (B).

to evaluate the relative extent by which Myo1 is likely to drive ring constriction by these two mechanisms. We consider the hypothetical case where Myo1 generates contractile force only through filament sliding without any impact on actin depolymerization. In this hypothetical cell the actin depolymerization-driven sliding (k) would be equal to that in *myo1 Δ tm*, but the activity of the motor (v) would be preserved as that in the WT cells. Simulation of the model based on this assumption gave the normalized contraction rate in this hypothetical cell α_h to be 0.0021 s^{-1} . Considering that the contribution of the motor to actin depolymerization accounts for a contraction rate of $\alpha_{WT} - \alpha_h$, and the total contribution of the motor to the contraction rate is $\alpha_{WT} - \alpha_{\Delta m}$, we estimate that 93% of the motor activity contributes to actin depolymerization, whereas only 7% contributes to direct filament sliding.

Contraction Rate Is Independent of the Initial Ring Size

Recent observations of actomyosin ring constriction in developing *C. elegans* early embryos revealed an intriguing invariant property, namely, the contraction time is independent of the cell size (Carvalho et al., 2009). In other words, the contraction rate (inverse of time in unit of s^{-1}) as we defined here is invariant with respect to different initial ring size. We note that in Carvalho et al. “rate” was used to denote contraction velocity measured in $\mu\text{m} \times \text{s}^{-1}$, which scaled with cell size. The authors proposed a mechanism of “structural memory,” referring to a certain precisely defined and spaced contractile unit shortening at a constant rate, the number of which scales with cell size. It was speculated that the scaling of contraction velocity with cell size may be an important property of rapidly dividing early embryos where the egg cytoplasm is successively divided during cell cycles of similar time. To test if this phenomenon

can be recapitulated by our model, we simulated rings of various initial sizes where larger rings on average contained proportionally large numbers of actin filaments with the same length and overlap distributions. This simulation found that, indeed, the contraction velocity scales with the initial ring size such that the duration of contraction is invariant (Figure 7A).

To test the aforementioned prediction, we investigated the relationship between the contraction rate and initial ring size at the bud neck. Because the natural variation of bud neck size is limited in a WT haploid population, we introduced the fluorescent probes into totally congenic diploid (2N), triploid (3N), and tetraploid (4N) strains (Pavelka et al., 2010) because cells of higher ploidy are in general larger in size, including the size of the bud neck, than cells of lower ploidy (Figures 7B–7F). Time-lapse measurements of cells ranging from 1N to 4N found that both the contraction rate α and the density loss rate β are invariant across cells with different bud neck diameters (Figures 7E and 7F). This finding confirmed our model prediction and suggests that the observed invariant contraction rate (or time) regardless of cell size is an intrinsic property of the actomyosin ring even in a unicellular organism. Interestingly, the initial density of the actin ring was also invariant across the range of bud neck sizes (Figures 7C and 7D), which would be predicted by invariant L_0 and S_0 .

DISCUSSION

The Function of Myosin II during Budding Yeast Cytokinesis

The role of the actomyosin ring and the mechanism by which force is generated to drive its constriction in dividing budding yeast cells have been a subject of considerable debate. The

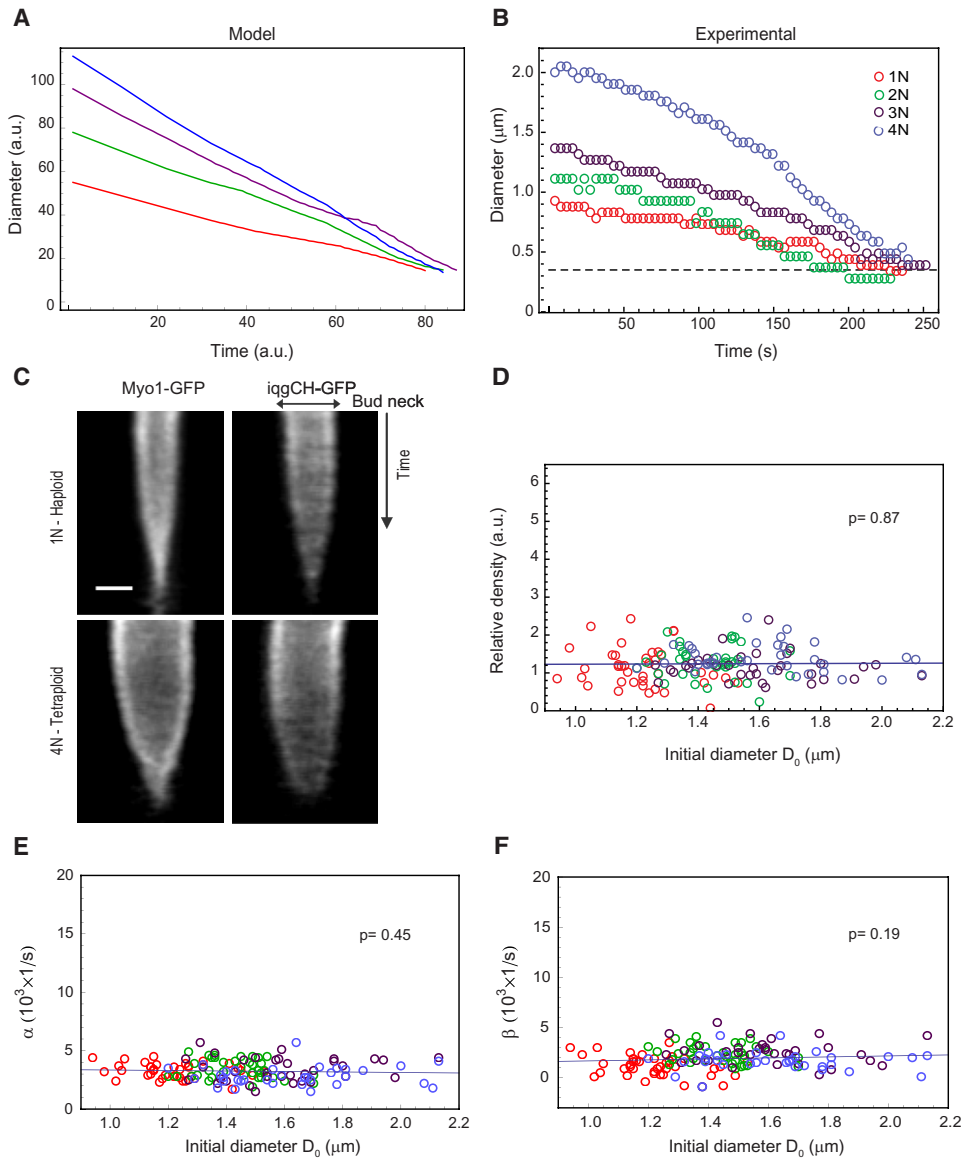


Figure 7. Actin Density and Contraction and Density Loss Rates Are Independent of the Initial Ring Size

(A) Model simulation of the ring diameter dynamics for ζ (0.37) and γ (0.27) for the WT cells with different initial ring diameter showing the contraction time to be independent of the initial ring size.

(B) Representative experimentally observed plots of ring diameter change over time for haploid (red), diploid (green), triploid (purple), and tetraploid (blue) cells. The dashed line shows the minimal size of the diameter below which the measurements are unreliable (see [Experimental Procedures](#) and [Figure S1](#)).

(C) Kymographs showing Myo1 and actin ring contraction and protein density dynamics for small haploid cells (first row) and large tetraploid cells (second row). The scale bar is 1 μm . The total time of each kymograph is 424 s.

(D–F) Relative actin density (D, see [Experimental Procedures](#)), contraction rate α (E), and actin density loss rate β (F) at the neck measured in cells with ploidy range from 1N to 4N (color coding as in B), showing independence of these parameters on the initial ring diameter. The p value is for comparison of the slope of the linear regression to a zero slope.

confusion largely stemmed from inconsistent phenotypes of *MYO1* gene knockout reported by different labs (Bi et al., 1998; Rancati et al., 2008; Tolliday et al., 2003; Watts et al., 1987) and the observation that Myo1 motor domain is not required for ring constriction. In the commonly used yeast strain backgrounds S288c and W303a, *MYO1* deletion leads to a nearly complete cytokinesis failure and cell lethality in the euploid genomic background. In S288c, viable $\Delta myo1$ spores with

massive growth and cytokinesis defects do occur at low frequencies (Rancati et al., 2008; Tolliday et al., 2003); however, upon further passage, innovation of novel cytokinetic mechanisms, not associated with formation of an actin ring, can be accomplished through aneuploidy (Rancati et al., 2008). Thus, in the experimental model used in our studies (S288c), Myo1 is indeed required for cytokinesis. However, this seems paradoxical with the observation that motor domain deletion could

largely rescue the cytokinesis defects of $\Delta myo1$ and the motorless ring can still constrict (Fang et al., 2010; Lord et al., 2005).

The results in this study help to resolve the aforementioned paradox by demonstrating that actin depolymerization, which occurs in the absence of the Myo1 motor activity, albeit at a slower rate, is a main mechanism for actomyosin ring constriction. Our experimental data and model analysis further indicate that the myosin II motor plays a role in actin depolymerization, although it is presently unclear whether this role could be direct, given that cofilin and myosin II binding to actin filaments are mutually exclusive (Galkin et al., 2011). A previous study in mammalian cultured cells also reported that inhibition of myosin II ATPase activity with blebbistatin prevented actin turnover (Guha et al., 2005). A role for myosin II in actin network disassembly was demonstrated recently in migrating fish keratocytes (Wilson et al., 2010). Another recent study suggests that tail retraction of migrating fibroblasts is a result of actin depolymerization, and myosin II facilitates this process by influencing filament alignment (Mseka and Cramer, 2011). Thus, promoting actin depolymerization may be a key function of the nonmuscle myosin II.

Force Generation via Actin Depolymerization

Force generation by depolymerizing cytoskeletal polymers is well known for the microtubule system. For example, microtubule depolymerization at the kinetochore contributes the force for poleward chromosome movement during anaphase (Bouck et al., 2008; Grishchuk et al., 2005; Koshland et al., 1988; Oguchi et al., 2011). During this process the linkage between kinetochore and the shrinking microtubule ends must be dynamically maintained. Actin polymerization and depolymerization have mostly been considered for the generation of protrusive forces via elongation at the barbed ends of filaments, whereas pointed end dynamics have mainly been implicated in actin subunit recycling. The possibility of actin depolymerization to produce contractile stress was recently proposed in a theoretical paper (Zumdieck et al., 2007) and discussed in Sun et al. (2010). Our observed effect of inhibition of actin depolymerization on the contraction rate provides the functional support for the idea of contractile force generation through actin depolymerization in the actomyosin ring.

Cofilin (Cof1) is the only essential actin-depolymerizing/severing protein in budding yeast and has been shown to promote rapid actin turnover in cortical actin patches and actin cables (De La Cruz, 2009; Fan et al., 2008; Lappalainen and Drubin, 1997; Moon et al., 1993). Our results show that Cof1 also regulates actin depolymerization in the contractile ring and plays an important role in ring contraction. Although several previous studies have clearly implicated ADF/cofilin in actomyosin ring assembly and/or contraction (Hotulainen et al., 2005; Kaji et al., 2003; Nakano and Mabuchi, 2006; Theriot, 1997), the precise functional consequence of cofilin-mediated actin depolymerization in cytokinesis has been unclear. Our analysis suggests that cofilin can be a key component of the force-generating machinery during cytokinesis. The action of cofilin must be augmented by actin crosslinking in order to generate contractile stress (Zumdieck et al., 2007). Our model suggests that such crosslinker does not need to track depolymerizing actin ends, as long as there is some probability (γ) for reestab-

lishing the crosslinking after filament shortening (Figure 5A). Iqg1 is a potential candidate for the critical actin crosslinker in our model because it is an actin crosslinker essential for cytokinesis (Epp and Chant, 1997; Lippincott and Li, 1998b). The constant value of γ assumed in the model implies that the crosslinker concentration in the ring should be constant during constriction, which was indeed observed for Iqg1 (data not shown).

A Quantitative Microscopic Model for Depolymerization-Driven Actomyosin Ring Contraction

The study of actomyosin-based contractile structures in different systems has led to models that consider different types of organization of actin filaments undergoing myosin motor-driven sliding (Carlsson, 2006; Carvalho et al., 2009; Stachowiak and O'Shaughnessy, 2009; Zemel and Mogilner, 2009; Zumdieck et al., 2007). The model proposed by Carlsson (2006) relies on actin turnover to ensure that the action of the motor leads to contraction and predicts that the contraction rate is limited by the rate of actin treadmilling. Another model by Stachowiak and O'Shaughnessy (2009) is characterized by a specific organization of stress fibers in which active myosin is placed with constant spacing, and the polarity of actin filaments alternates periodically along the fiber axis. Both these models assume the myosin motor being the necessary contractile stress generator. The model of Zumdieck et al. (2007) illuminates the possibility of actin depolymerization coupled with filament crosslinking to constitute an alternative driving force for filament sliding. The actin filaments are assumed to undergo treadmilling, which may be consistent with the observation made in fission yeast by Pelham and Chang (2002).

The bottom-up model, based on microscopic elements of dynamic actin filaments, actin crosslinker, and myosin II motor, allowed us to compute macroscopic parameters describing contraction dynamics to compare with the experimentally measured values. An attractive feature of our model is the lack of requirement for precise filament configuration. Three-dimensional electron microscopy reconstruction of the fission yeast contractile ring at the start or during contraction showed a mixed filament polarity with perhaps a slight bias in one orientation (Kamasaki et al., 2007). The model predicts that indeed no bias in filament orientation is required for the ring contraction (Figure 6), but a small nonzero ζ is needed to obtain the observed rate of the actin loss in the contracting ring.

Our model correctly predicts the independence of contraction rate (or time) on the initial size of the actomyosin ring. This phenomenon was recently reported for early *C. elegans* embryos undergoing rapid mitotic cleavages (Carvalho et al., 2009). The observation in yeast suggests that this property is a predictable outcome of an evolutionarily conserved aspect of actomyosin ring mechanics rather than an evolutionary innovation to accomplish precisely timed, consecutive embryonic divisions. In contrast to the model proposed by Carvalho et al. (2009), our model does not require spatially separate structural units of a fixed size, but rather, actin filaments can vary randomly in orientation, length, and the amount of overlap within the defined distributions. Having overlapping filaments throughout the ring is likely to be important for maintaining its structural integrity during contraction and is also consistent with the electron microscopy

reconstruction of the fission yeast ring (Kamasaki et al., 2007). Finally, although our model simulation predicts that in budding yeast actin depolymerization is the main mechanism driving actomyosin ring constriction, the model itself does not exclude a direct role for the motor and may thus be generalizable toward systems where filament sliding is more predominantly driven by myosin II.

EXPERIMENTAL PROCEDURES

Yeast Strains and Plasmids

Media and genetic techniques were as previously described (Sherman et al., 1974). Yeast strains and plasmid construction are in Supplemental Experimental Procedures).

Live Imaging and Image Analysis

Three-dimensional confocal imaging was accomplished using a Carl Zeiss 510 Meta laser-scanning microscope equipped with a Confocor 3 detection unit. This detection unit employs single-photon counting avalanche photodiodes that are crucial for highly sensitive and quantitative imaging of contractile ring dynamics in yeast cells with native expression levels. For this study, two general methods were employed. For time-lapse imaging to observe density changes, it was crucial to avoid photobleaching. Therefore, the pinhole was set completely open with relatively low-laser power, and a Carl Zeiss C-Apochromat 40x 1.2 NA water objective was used. Under this condition, no photobleaching was observed in our time-lapse movies.

All image processing was done using ImageJ, a free open source Java program from the NIH combined with several custom plug-ins written in-house. For diameter measurements a thick line was drawn through the bud neck parallel to the ring such that the thickness of the line encompassed the observed thickness of the ring (Figure S1A). A summed and smoothed profile of this region as a function of time represents a kymograph from which the background intensity was subtracted. For each time point profile, the maximum intensity value was measured. Then the diameter defined as the profile width at the half-maximum level was computed by searching inward from each edge of the profile and then interpolating near the half-maximum (Figure S1B). Because our measurements were near the limit of optical resolution, we measured the distance between the positions of the profile peaks representing the initial ring diameter and compared it to the half-maximum width measurement. We found that the difference was roughly 350 nm. The end time point of the contraction process for each individual cell measurement was defined as the time when the measured diameter reaches this critical value (see Figure S1B). The start point of the contraction phase was determined as the start point of the diameter curve's linear segment (see Figure 2A), which often correlates in time to a sudden inward movement of the spindle pole bodies, labeled with Spc42-mCherry, an event caused by spindle breaking down at mitotic exit. To generate the contraction curves, each diameter measurement was normalized to the initial diameter determined as the average of the ten time points before the start of contraction. The normalized contraction rate was computed in *Mathematica* (Wolfram Research, Champaign, IL, USA) as the slope of a linear fitting of the diameter profile during the defined contraction window. All images for presentation were adjusted for contrast to clearly show the maximum and minimum intensities. See Supplemental Experimental Procedures for the methods for actin density measurements.

Estimate of α and β

To estimate α and β from experimental measurements of ring diameter and intensity changes over contraction time (e.g., plots in Figure 2A), we used the following approximation:

$$D(t) = D(0)\exp(-\alpha \cdot t) \approx D(0)(1 - \alpha \cdot t),$$

where the last expression is obtained as the leading one in expansion when α is small. If α is constant, this approximation leads to the constant speed of contraction (measured in $\mu\text{m/s}$), as observed (Figure 2A), which is equal to $\alpha D(0)$ and scales with the initial ring diameter. Similar reasoning leads to approximated protein density dynamics $c(t) = c(0)\exp(-\beta \cdot t) \approx c(0)(1 - \beta \cdot t)$.

Again, it is important to remind that α and β have the dimension of s^{-1} , not $\mu\text{m/s}$.

Statistical Analysis

Statistical analysis of the data was performed in *Mathematica* and Excel. The p values were determined by the Student's t test.

Model Computation

Model simulations were made using custom code written in *Mathematica*. The initial orientation, size, and position of actin filaments were generated randomly using built-in *Mathematica* random number generator. The search for model parameters, such as the clustering coefficient ζ and the frequency γ of productive actin depolymerization, was made by running, for each parameter set in the range $-1 < \zeta < 1$ with the step 0.1 and $0 < \gamma < 1$ with the step 0.05, 200 simulations starting with random initial filament structures, computing the average values of contraction rate α and actin loss rate β . The values of other parameters were selected as follows: initial filament mean size and SD were $L_0 = 1$ and $\delta L = 0.15$, respectively; the overlapping mean size and SD were $S_0 = 0.75$ and $\delta S = 0.25$, respectively; and initial number of filaments was 20. The depolymerization rate was chosen to be equal to $k = 0.01$ and the myosin motor speed $v = 0.0025$ for the WT simulation. For the motorless simulation the depolymerization rate was reduced to $k = 0.0061$, and myosin motor speed v was set to zero. To obtain a finer resolution in the subregion in the parameter space in the range $0.2 < \zeta$, $\gamma < 0.4$, similar computations were done 300 times for each parameter set with step size of 0.02 in both parameters. In the simulation for the case of *cof1-22* mutant, the values $k = 0.0076$ and $v = 0.0025$ were used. For the simulations for different initial neck size, 10, 15, 20, and 25 were used as the initial number of filaments.

SUPPLEMENTAL INFORMATION

Supplemental Information includes two figures, six movies, and Supplemental Experimental Procedures and can be found with this article online at doi:10.1016/j.devcel.2012.04.015.

ACKNOWLEDGMENTS

The authors thank J. Zhu (Stowers Institute) for assistance in karyotyping, A. Mushagian (Stowers Institute) for supervision of A.K., and E. de Melo (ITQB and New University of Lisbon, Portugal), M. de Sousa, F. Carneiro, and J. C. Lopes (GABBA program, Faculty of Medicine, University of Porto, Portugal) for their mentoring to I.M.P. The authors thank the reviewers of this manuscript for the constructive and helpful comments. This work was supported by NIH grant RO1GM059964 to R.L.

Received: August 23, 2011

Revised: January 27, 2012

Accepted: April 19, 2012

Published online: June 11, 2012

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